## **Droplet** CFD

n Energent•s Variable Phase Turbine [1-2] (VPT) the "uid at the inlet is liquid, "ashes inside the nozzle upstream of the turbine rotor, and is two-phase inside the rotor blade passage. A previous article [3] discussed calculating the trajectories of droplets inside the turbine rotor.

In the converging section of the nozzle, the pressure decreases. When it declines to the saturation pressure, vapor bubbles form. At this pressure, the liquid is the continuous phase, the vapor the dispersed phase. ssureuvec i.rvtaletthis pressure,

the calculations to develop a reduced order model that can be incorporated into traditional CFD codes and 1-D nozzle codes. Experimental work has been found for model problems to begin investigating computationally. By "nding problems to study that have been investigated experimentally, the methodology used in the CFD simulations can be validated.

A starting point is to examine the "ow "eld around a single liquid





Figure 7 shows the re"ected and transmitted shock waves, well as the unsteady "ow conditions both inside and behind t cylinder cloud.



Figure 4 Drag coef"cient non-dimensionalized using the deforming column frontal diameter, for a threshold value of 0.95 of the liquid volume fraction, for several different incident Mach numbers.

## By using the actual column diameter instead of the initial diameter

the drag coef"cient shows signi"cantly less variation during the breakup period that is simulated, Figure 4.

In the application of interest, the droplet is not in isolation, but is part of a cloud. At Sandia National Laboratory [5], experimentation at test impacts a conducted on a planar shock wave impacting allowing for the view of solid particles.

coupling terms are still considered. The 1-D model equations the simulation focuses on the early stage of the experime blved do not include the unclosed "uctuation terms creat when the particles have not yet moved and can be assumed in blue volume-averaging procedure, such as the Reyno "xed in space. The 3-D particle cloud is modeled by an effects." This is a reasonable assumption in dilute multiphase "ov staggered cylinders, Figure . With the stagger arrangement were fr, in dense "ows this assumption may not be appropriate the open cross sectional area varies by less than 1.5%, Figure 6.

The volume fraction is nearly constant through the curtaine friescellaneous particle forces are assumed to be included in this 2-D model, the Euler equations are solved. The nutragicaefficient for the quasi-steady drag force on a single particle method implicitly contains numerical viscosity.



Figure 5 Array of staggered cylinders.

 $F_{i}^{qs} = \frac{1}{2} C_{D} A_{p} | u_{i} - v_{i} | (u_{i} - v_{i})$ 

where  $A_p$  is the particle cross-sect  $Q_p$  is the drag coef"cient, and u and v are the continuous and dispersed phase velocities respectively. For the time period considered the particle is "xec space,  $s_{ij} = 0$ . For the 2-D particle, the cross-section area is its diamete  $D_p$ . The drag coef"cie  $Q_{ij}$  was determined by "nding the value that best matches the re"ected and transmitted sho locations and magnitudes of the 2-D solution.

Figure 8 - Figure 10 compare the solution of the 1-D model w a planar average of the 2-D result at the non-dimensionalize time of 3.5. The particle curtain is located between -0.5 < x -0.5. For the plots of density and velocity, the 2-D results apper to oscillate around the 1-D model results for a signi"cant portion of the solution. For these pro"les, two additional cases are show where the drag coef"cient is increased and decreased by 30° Small, yet noticeable, differences can be observed in the sho locations. This suggests that the methodology used is adequate evaluate an overall mean drag coef"cient.

continued on page 8

7





Figure 8 Comparison of the density from the 1-D model and with the planar average of the 2-D model at t=3.5. In addition for the 1-D model, the drag coef"cient was varied by +/- 30%.

Figure 9 Comparison of the velocity from the 1-D model with the planar average of the 2-D model at t=3.5. In addition for the 1-D model, the drag coef"cient was varied by +/- 30%.

In Figure 10 the planar averaged pressure in the 2-D result is consistently lower than that predicted by the 1-D model inside the particle cloud and downstream of the trailing edge until x 1.5. This is attributed to the "uctuations associated with the vortical structures (see Figure 7), which is a behavior that the 1-D model,

